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# **The** TRANSOPTR **Model of the ISAC Drift Tube Linear Accelerator - Part I: Longitudinal Verification**

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**Abstract:** This note records the initial implementation of the ISAC Drift Tube Linac in the envelope code TRANSOPTR. As part of an effort to explore and map the vast set of longitudinal DTL tunes, a full model has been built using synthetic electric field distributions, generated in the Poisson solver Opera-2D, from detailed technical schematics of the as-built interdigital structure for each cavity of the accelerator. Exploiting TRANSOPTR's subsecond execution times means surveys of accelerated DTL output beam parameters versus cavity  $(V_s, \phi)$  can be produced in a matter of minutes. Model computed output energy-phase curves at constant electric field amplitude are compared to on-line measurements using an oxygen beam, as part of a verification of model's longitudinal predictive power.

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# **Contents**



<span id="page-2-0"></span>This document records the TRANSOPTR implementation of the ISAC-DTL [\[1\]](#page-34-1), a variable energy postaccelerator for low intensity radioisotope beams (RIB), whose design operating tank and buncher energies are listed in Table [1.](#page-2-1) Users of DTL beams are high energy nuclear physics experiments who by their nature constantly request different beam energies. This is achieved by sequential tuning of each tank and buncher to its design energy, up until the requested output energy is reached, with downstream tanks powered off. For example, a desired final output energy of 0.9 MeV/u would be achieved by setting up to Buncher-3 at design, with Tank-4 then being tuned to produce the desired output and Tank-5 off. It is important to stress that for the DTL, design energy is not the highest attainable tank energy. Instead, the design velocity for each tank produces a minimized increase in energy spread and bunch length, over the length of the machine.



<span id="page-2-1"></span>Table 1: Design output energies for the ISAC-DTL [\[2\]](#page-34-2). Note Tank-1 in practice is operated at 0.238 MeV/u while the reference quotes 0.236 MeV/u.

The mass-to-charge ratio for DTL accelerated beams by design is  $2 < A/q < 6$ . From electrodynamics, it follows that beams with the same A/q should have the same set of accelerating gradients  $(V_s)$  [\[3\]](#page-34-3) and phases. For example,  $^{20}Ne^{5+}$  and  $^{16}O^{4+}$  should both have the same voltage gradient setpoints, while  $^{20}$ Ne<sup>5+</sup> and  $^{12}$ C<sup>3+</sup> will have gradients differing by a factor of 4/5, owing to the A/q difference. However, in terms of cavity phasing, there should be no difference since there is a unique time-of-flight profile between drift tubes and the injection beam velocity is unchanging.

Output DTL energy is measured using a  $90^\circ$  analyzing magnet coupled with a wire scanner (*the Harp*) [\[4\]](#page-34-4), allowing operators to view the energy spectrum of the output bunch distribution. Operators power on the RF amplifiers at low amplitude to limit multipactoring behavior, which also produces minimal initial tank acceleration upon powerup. Tank setup or energy ramping proceeds when an operator manually tunes the parameter  $V_s$  while the energy spectrum is monitored on the Harp.

# **DRIFT TUBE LINAC**



<span id="page-3-0"></span>Figure 1: Overview of the ISAC Drift Tube Linac, an 106.08 MHz, separated function, interdigital Hmode heavy-ion linac. The quadrupole triplet assemblies, which provide transverse focusing, are interspersed between the tanks along the structure. Cutaway of tanks shows internal structure, where the IH drift tube supports can be seen attached to resonators at the top and bottom of each tank. The three bunchers are split-ring three-gap structures. Each buncher cavity is identical in dimension, while the gaps differ. Beam propagation occurs from left to right. Image courtesy M. Marchetto.

During the energy ramp, the RF phase is scanned to minimize the output beam energy spread [\[5\]](#page-34-5). This phase optimization is necessary to maintain beam transmission and keep the energy distribution visible on the Harp. Simultaneously, operators are also required to re-tune the quadrupole triplets and scale the HEBT quadrupole currents for the increased beam energy, not to mention continually tune the  $90^\circ$  magnet's field to keep beam visible on the diagnostic. The final DTL beam energy is computed from the current necessary to deflect beam at the HEBT1 station, using a BI parametrization of the magnet.

The above mentioned procedure is both time consuming and demanding on operators, due to the number of tasks that must be carried out simultaneously. In practice, it is not uncommon for an entire 8-hour shift to be consumed ramping the DTL to its full design energy. Further, the heavy procedural overhead produces a system which is prone to human error or oversight. The end result is that tuning the DTL is considered one of the most complex tasks to be carried out by operators.

However, from a physical standpoint, the intense procedural load is difficult to justify given the extreme reproducibility of the linac, in addition to the set of configurations for which the DTL is used in ISAC beam delivery. One would be justified in arguing that *a priori*, since it always ends up tuned within its nominal performance envelope and with a set of unchanging design energies, much of the initial setup procedure can be considerably reduced, by obtaining a full mapping of the energy-phase relation for each of the tanks and bunchers of the linac. This would allow operators to skip much of the ramping procedure, and go directly to the desired energy. Even better, if one can quickly model this, in a manner that agrees with the machine's behaviour, then much of this procedure can be automated and offloaded from the operators' plate altogether.

To date, the project to implement a full end-to-end model of the ISAC linac in TRANSOPTR has scrutinized OLIS and ILT [\[6,](#page-34-6) [7\]](#page-34-7), added and benchmarked an optr-RFQ capability for the ISAC-RFQ [\[8\]](#page-34-8) and also built and benchmarked the medium, high and superconducting energy beamlines [\[9,](#page-34-9) [10,](#page-34-10) [11\]](#page-35-0). An ISAC-II model of SCRF is, in parallel, underway [\[12,](#page-35-1) [13\]](#page-35-2). A tracking of global time was implemented in optr, to allow faithful modelling of inter-cavity time of flight effects [\[14\]](#page-35-3) in all linac simulations.

Last year, an investigation of the DTL design tune found anomalies in quadrupole BI parametrizations and succeeded at achieving a high transmission on-line tune at tank-2 energy using model computed quadrupole setpoints. The Langevin-like BI parametriza-tions from that report [\[15\]](#page-35-4) are used in the model.

And so, this document represents the laying of a cornerstone for the full TRANSOPTR model of the ISAC accelerator and with it the possibility of further developing the methods and principles of ISAC beam delivery. As the present DTL implementation is adaptable to the nascent HLA infrastructure at TRIUMF [\[16\]](#page-35-5), the significance of this model and its capabilities should be evident.

# <span id="page-4-0"></span>**2 On-Line DTL Energy-Phase Measurements**

A comprehensive beam based measurement of the energy-phase output of each DTL tank and buncher was organized by the author in December 2017, using a  ${}^{16}O^{4+}$  beam from the microwave source. Over a continuous 48 hour period, a team of volunteers assisted in performing these measurements, which provided a detailed mapping of output cavity energy versus cavity tuneable parameters  $(V_s, \phi)$ . The run plan and data acquisition protocol for this experiment are shown in Appendix [A.](#page-36-1) The data, shown in Figures [2](#page-5-0) and [3,](#page-6-0) preceded construction of the model and was consequently catalogued and stored until the present, where it is to be used for model verification, in Section [5.1.](#page-19-1)



<span id="page-5-0"></span>Figure 2: Energy-Phase lines for the MEBT Rebuncher, which precedes the ISAC-DTL, and of its bunchers, as measured with  $16O^{4+}$ . Amplitude settings, as set by operators via EPICS, are shown in the legend for each plot. Each line shows the linac output for constant amplitude, implying constant on-axis electric field. The MEBT Rebuncher was mapped and is shown, though is not used in the present analysis.



<span id="page-6-0"></span>Figure 3: Energy-Phase lines for the ISAC-DTL, as measured with  $^{16}O^{4+}$ . Amplitude settings, as set by operators via EPICS, are shown in the legend for each plot. Each line shows the linac output for constant amplitude, implying constant on-axis electric field.

# <span id="page-7-0"></span>**3** Opera-2D **Longitudinal Electric Field Model**

In [\[3\]](#page-34-3), Baartman develops the F-matrix formalism for an axially symmetric, time varying electric field, as found in many linacs, including the DTL. While the assumption of full axial symmetry is not perfectly true, owing to the presence of a small dipole component arising from the IH structure's stems, the tank design aims to minimize this, and for the ISAC-DTL it is negligible [\[17\]](#page-35-6). To produce the TRANSOPTR model, the dimensions of the as-built drift tubes were kindly provided by R. Laxdal and B. Waraich and were used to cross-reference technical drawings obtained from the TRIUMF Design Office. As the acceleration through the linac depends entirely on the field  $\mathcal{E}(s)$ , incorrect dimensions would likely produce an output energy spectrum which would disagree with observations.

### <span id="page-7-1"></span>**3.1** Opera-2D E(s) **Model and Dimensions**

The electromagnetic code Opera-2D was used to generate a model of the drift tubes for each resonator in the DTL using the provided dimensions, following a geometric model shown in Figure [4.](#page-7-2) The dimensions for the bunchers were obtained from [\[18\]](#page-35-7), which records the design dimensions from the team at INR-RAS in Moscow. All dimensions and corre-sponding drawings are listed in Tables [2,](#page-8-0) [3](#page-9-0) and [4.](#page-10-0)



<span id="page-7-2"></span>Figure 4: Geometric representation of the Opera-2D model used for the axially symmetric IHstructure geometry, showing a single drift tube as an example, between two outer tubes connected to the cavity wall. The axis of radial symmetry is shown in purple, labelled  $S_r$ . The outer tubes are grounded (GND) while the inner tube(s) alternate between a constant value of  $\pm V$ . Tube dimensions, including rounded edge radii and tube lengths have been extracted from design drawings or specified reference.



Table 2: Physical drift tube dimensions for (**Left**) DTL Tank-1, extracted from drawing IRF0571D.dwg (Fig. [5\)](#page-8-1) and (**Right**) DTL Tank-2, from drawing IRF1102D.dwg (Fig. [6\)](#page-9-1), obtained from the TRIUMF Design Office.

<span id="page-8-0"></span>

<span id="page-8-1"></span>Figure 5: **Top:** IRF0571D.dwg, showing the interdigital structure within DTL Tank-1. Extracted tube lengths are listed in Table [2.](#page-8-0)



<span id="page-9-1"></span>Figure 6: IRF1102D.dwg, showing the interdigital structure within DTL Tank-2. Extracted tube lengths are listed in Table [2.](#page-8-0)

tube start [in]	tube end [in]	v [kv]	tube start [in]	tube end [in]	V [kV]
0.0000	1.5230	0.0	0.0000	1.4170	0.0
2.1900	3.2850	$-1.0$	2.3450	3.7250	$-1.0$
3.9620	5.0630	1.0	4.6650	6.0490	1.0
5.7910	6.8750	$-1.0$	7.0410	8.4060	$-1.0$
7.6630	8.7350	1.0	9.4570	10.8100	1.0
9.5650	10.6380	$-1.0$	11.9020	13.2540	$-1.0$
11.5040	12.5840	1.0	14.3180	15.7460	1.0
13.4810	14.5760	$-1.0$	16.8900	18.2780	$-1.0$
15.4930	16.6160	1.0	19.4340	20.8560	1.0
17.5380	18.6980	$-1.0$	22.0080	23.4810	$-1.0$
19.6220	20.8250	1.0	24.6100	26.1450	1.0
21.7410	23.0020	$-1.0$	27.2470	28.8510	$-1.0$
23.8910	25,2240	1.0	29.9130	31.5790	1.0
26.0730	27.4830	$-1.0$	32.6280	34.3290	$-1.0$
28.2950	29.7540	1.0	35,3870	37.3130	0.0
30.5750	32.1950	0.0			

<span id="page-9-0"></span>Table 3: Physical drift tube dimensions for (**Left**) DTL Tank-3, extracted from drawing IRF1183D.dwg (Fig. [7\)](#page-10-1) and (**Right**) DTL Tank-4, from drawing IRF1243D.dwg (Fig. [8,](#page-11-0) top), obtained from the TRIUMF Design Office.

tube start [in]	tube end [in]	V [kV]
0.0000	1.4490	0.0
2.6850	4.2220	$-1.0$
5.4700	7.0240	1.0
8.2970	9.8600	$-1.0$
11.1640	12.7440	1.0
14.0580	15.6660	$-1.0$
16.9860	18.6280	1.0
19.9860	21.6340	$-1.0$
22.9340	24.6790	1.0
25.9510	27.7650	$-1.0$
28.9960	30.8960	1.0
32.0610	34.0360	$-1.0$
35.1820	37.1940	1.0
38.3490	40.4630	0.0

<span id="page-10-0"></span>Table 4: Physical dimensions extracted from drawing IRF1302D.dwg (Fig. [8,](#page-11-0) bottom) for ISAC-DTL Tank-5, 106.08 MHz.



<span id="page-10-1"></span>Figure 7: **Top:** IRF1183D.dwg, showing the interdigital structure within DTL Tank-3. Extracted tube lengths are listed in Table [3.](#page-9-0)





(57.195)

<span id="page-11-0"></span>Figure 8: **Top:** IRF1243D.dwg, showing the interdigital structure within DTL Tank-4. **Bottom:** IRF1303D.dwg, DTL Tank-5. Extracted tube lengths for both tanks are listed in Table [3.](#page-9-0) **Drawings not to scale with respect to each other.**

Additional simulation parameters for the DTL tanks, corresponding to those shown in Figure [4](#page-7-2) are shown in Table [5.](#page-12-0) For the bunchers, the extracted dimensions are shown in Table [6,](#page-12-1) with the tube-specific parameters listed in Table [7.](#page-13-0) The technical drawing IRF0965D.dwg is shown in Figure [9,](#page-13-1) obtained from INR-RAS, representing the thre-gap DTL bunchers.<sup>[1](#page-0-0)</sup>. Note that the tube dimensions shown in the figure are a placeholder, with the final as-built dimensions listed in ref. [\[18\]](#page-35-7).

Cavity	No. Tubes	$W_t$ [in]	$r_b = r_t$ [in]	$r_i$ [in]	$s_{end}$ [in]
IH Tank-1	8	0.1970	0.0938	0.2755	12.8923
IH Tank-2	12	0.2953	0.0985	0.19685	21.6850
IH Tank-3	14	0.1970	0.0938	0.3150	32.1950
IH Tank-4	13	0.1970	0.0938	0.3150	37.3130
IH Tank-5	12.	0.1970	0.0938	0.3150	40.4630

<span id="page-12-0"></span>Table 5: Extracted tube dimensions for each tank, showing the cavity name, number of drift tubes (excluding grounded end tubes), the tube thickness  $W_t$ , rounded edge radii, which is equal for the IH cavities, in addition to the tube inner radius  $r_i$  and end coordinate of the field map  $\mathcal{E}(s)$ . All dimensions are in inches as presented on the technical drawings.



<span id="page-12-1"></span>Table 6: Physical drift tube dimensions for (**Top-Left**) DTL Buncher-1, (**Top-Right**) DTL Buncher-2 and (**Bottom**) DTL Buncher-3, all obtained from [\[18\]](#page-35-7).

<sup>&</sup>lt;sup>1</sup>Note the transliteration errors from cyrillic to latin characters in Figure [9](#page-13-1)



<span id="page-13-1"></span>Figure 9: Digitized design drawing for the 3-gap spiral DTL bunchers, originally from INR-RAS, used for Opera2D field map generation. Note the transliteration errors representing original cyrillic characters. Dimensions as shown in the drawing are in millimeters, and are valid. Obtained from the TRIUMF Design Office. Note, the technical drawing is identical to that for the MEBT bunch rotator, shown in [\[9\]](#page-34-9).

Cavity	No. Tubes $ W_t$ [in] $ r_b$ [in] $ r_t$ [in] $ r_i$ [in] $ s_{end}$ [in]			
Buncher-1	$\overline{2}$	$\vert$ 0.3150 $\vert$ 0.0987 $\vert$ 0.1392 $\vert$ 0.4918 $\vert$ 4.8226		
Buncher-2	$\mathcal{P}$	$\vert$ 0.3150 $\vert$ 0.0987 $\vert$ 0.1392 $\vert$ 0.4918 $\vert$ 5.8466		
Buncher-3	2	$\vert$ 0.3150 $\vert$ 0.0987 $\vert$ 0.1392 $\vert$ 0.4918 $\vert$ 6.6405		

<span id="page-13-0"></span>Table 7: Extracted tube dimensions for each buncher, showing the cavity name, number of drift tubes (excluding grounded end tubes), the tube thickness  $W_t$ , rounded edge radii  $r_t$  and  $r_b$ , in addition to the tube inner radius  $r_i$  and end coordinate of the field map  $\mathcal{E}(s)$ . All dimensions are in inches as presented on the technical drawings.

### <span id="page-14-0"></span>**3.2** Opera-2D **On-Axis Electric Fields**

The model for each tank and buncher was processed in Opera-2D allowing for the generation of an on-axis electric field distribution  $\mathcal{E}(s)$  for each RF component of the DTL, enabling the use of TRANSOPTR subroutine linac for each cavity, with outputs listed in Table [8.](#page-14-1)

It is noted that unlike published bead pull measurements (e.g. Fig. 3 in [\[19\]](#page-35-8)), the Opera-2D output field distributions possess varying peak heights along the length  $s$  of the structures and that in this implementation, no attempt has been made to manipulate the tube voltages to flatten the peak height distribution. This is noted as a possible model tuning strategy, though will not be further explored in this report.



<span id="page-14-1"></span>Table 8: List of figures showing the Opera-2D computed fields associated with each RF cavity in the ISAC DTL. Figure placement specified as top (t), middle (m) or bottom (b).



<span id="page-15-0"></span>Figure 10: Normalized longitudinal electric field  $\mathcal{E}(s)$  for the ISAC Drift Tube Linac, showing (Top) Tank-1 and Tank-2 (**Bottom**)



<span id="page-16-0"></span>Figure 11: Normalized longitudinal electric field  $\mathcal{E}(s)$  for the ISAC Drift Tube Linac, showing Tank-3 (**Top**), Tank-4 (**Middle**) and Tank-5 (**Bottom**).



<span id="page-17-0"></span>Figure 12: Normalized longitudinal electric field  $\mathcal{E}(s)$  for the ISAC Drift Tube Linac, showing Buncher-1 (**Top**), Buncher-2 (Middle) and Buncher-3 (**Bottom**).

# <span id="page-18-0"></span>**4** /acc**-database** TRANSOPTR **Implementation**

The model for the linac uses technical drawing IRF1002D.dwg, shown in Figure [1,](#page-3-0) for tank, buncher and quad positioning along the optical axis s. The field maps from Section [3](#page-7-0) are centered at the geometric midpoint of each tank or buncher, referenced to the outer tank surfaces. The sequence dt1\_db0.xml in the /acc database has been defined from the above drawing and the sequence layout is shown in Table [9.](#page-18-1)



<span id="page-18-1"></span>Table 9: Sequence dt1\_db0: the ISAC Drift Tube Linac with reference start and end coordinates from which elements and their positions along the optical axis, S in millimeters were extracted. Quadrupole effective lengths obtained from [\[20\]](#page-35-9). The TRANSOPTR element (subroutine) type, in addition to the element length are also displayed.

The quadrupole triplet effective lengths were set to 5.8cm and 8.7cm for the outer and inner quadrupoles, respectively [\[20\]](#page-35-9). A detailed analysis of the BI relationship for the DTL quadrupoles may be found in [\[15\]](#page-35-4).

## <span id="page-19-0"></span>**5** TRANSOPTR **Simulations**

An initial comparison was performed with an  $A/\alpha=6$  beam using a starting tune based on parameters obtained from [\[2\]](#page-34-2) for the initial sigma matrix parameters. Optr was run in mode 5, featuring the addition of global time tracking [\[14\]](#page-35-3), which allows for a time-of-flight based relative tank phasing. The /acc database implementation of the HEBT line [\[10\]](#page-34-10) was used, along with dtl\_db0.xml to generate a TRANSOPTR system file representing the entire ISAC-DTL, starting at DTL:FC0 and terminating at the HEBT1 high energy diagnostic station (Prague).

The longitudinal tune of the linac was found by performing a comprehensive mapping of the  $(V<sub>s</sub>, \phi)$  configuration space of each tank and buncher. This data was used to find the opti-mum accelerating tank setpoints producing the design energy profile specified in Table [1,](#page-2-1) in Section [1.](#page-2-0) These also allowed for the generation of TRANSOPTR computed energy-phase curves at constant  $V_s$ , produced to have a maximum energy identical to each constant am-plitude line in the datasets from Section [2.](#page-4-0) The comparison of the optr scaling factor  $V_s$ with the EPICS PV setpoints allows for the extraction of an initial calibration between both, important for on-line use.

### <span id="page-19-1"></span>**5.1 Beam Based** TRANSOPTR **Amplitude Calibration**

The simulation data from the previous section produces a set of energy-phase curves which can serve as an observable prediction to be compared with beam-based measurements, an initial test of model validity. For this, a modified version of topology [\[21\]](#page-35-10), which found the minimum  $V_i$  that can produce an output energy  $E_i(V_i,\phi)$  corresponding to the peak energy of each beam-acquired dataset.

The phases which are presented are relative, no calibration has yet been established between the master RF clock which is controlled in EPICS and the TRANSOPTR model's phase. In principle, this means the value of the phase will disagree between model and machine, but the relative phase differences between cavities should be the same, taking into account time-of-flight. The specified tolerance between model maximum output energy and beambased data, for each scan, was  $\Delta E/E = 10^{-5}$ . This corresponds to finding the maximum  $\mathcal{T}(V_{s}, \phi_{0})$  [\[22\]](#page-35-11). For DTL Tank-1, Buncher-1 and Tank-2, this is shown in Figure [13.](#page-20-0) For Buncher-2, Tank-3 and Buncher-3, it is shown in Figure [14.](#page-21-0) Tanks 4 and 5 are in Figure [15.](#page-22-0) The figures show the correspondence established between the EPICS PV's from the beam measurements and the model. Initial conditions are shown in Tables [10](#page-22-1) and [11.](#page-22-2)



<span id="page-20-0"></span>Figure 13: Top-Left: Calibration of EPICS voltage scaling factor and TRANSOPTR V<sub>s</sub> for DTL Tank-1. **Top-Right:** corresponding TRANSOPTR  $E(\phi_0)$  curves. **Middle:** Corresponding curves for DTL Buncher-1. **Bottom:** Corresponding curves for DTL Tank-2. Based on <sup>16</sup>O<sup>4+</sup> data acquired at HEBT1 diagnostic station. Note the phase axes have not been calibrated in this figure.



<span id="page-21-0"></span>Figure 14: **Top-Left:** Calibration of EPICS voltage scaling factor and TRANSOPTR  $V_s$  for DTL Buncher-2. **Top-Right:** corresponding TRANSOPTR E( $\phi$ <sub>0</sub>) curves. **Middle:** Corresponding curves for DTL Tank-3. Bottom: Corresponding curves for DTL Buncher-3. Based on <sup>16</sup>O<sup>4+</sup> data acquired at HEBT1 diagnostic station. Note the phase axes have not been calibrated in this figure.



<span id="page-22-0"></span>Figure 15: Top-Left: Calibration of EPICS voltage scaling factor and TRANSOPTR V<sub>s</sub> for DTL Tank-4. **Top-Right:** corresponding TRANSOPTR  $E(\phi_0)$  curves. **Bottom:** Corresponding curves for DTL Tank-5. Based on  ${}^{16}O^{4+}$  data acquired at HEBT1 diagnostic station. Note the phase axes have not been calibrated in this figure.

$$
\frac{\epsilon_x \, [\mu \text{m}] \quad \epsilon_y \, [\mu \text{m}] \quad \epsilon_z \, [\mu \text{m}] \quad r_{xx'} \quad r_{yy'} \quad r_{zz'}}{39.8 \quad 34.0 \quad 26.7 \quad -0.89 \quad -0.90 \quad -0.62}
$$

<span id="page-22-1"></span>Table 10: Starting emittances and bunch correlation coefficients between canonical coordinates used for TRANSOPTR-DTL simulations.

			$E_i$ [MeV] $m_0$ [u] q [e] 2x-rms [cm] 2 $P_x$ -rms [mrad] 2y-rms [cm] 2 $P_y$ -rms [mrad] 2-zrms [cm] 2 $P_z$ -rms [mrad]				
4.59	30.0	0.36	23.60	0.34	22.54	0.14	25.17

<span id="page-22-2"></span>Table 11: TRANSOPTR DTL injection parameters at start of sequence dtl\_db0.xml, used for present analysis. The parameters are derived from [\[2\]](#page-34-2).



<span id="page-23-0"></span>Figure 16: Measured energy spread (FWHM) at HEBT1 HE diagnostic station for  $^{16}O^{4+}$  beam, showing normalized, operator estimated energy spread compared to value for maximum output energy at each amplitude, featuring minimized energy spread. The phases are measured with respect to the optimum accelerating phase, showing the rapid growth in energy spread with positive or negative dephasing.

Broadly speaking, the agreement for the cases shown in Figures [13,](#page-20-0) [14](#page-21-0) and [15](#page-22-0) can be described as good, with minor disagreements between model and beam. However, for Tank-3 and Buncher-3, in Figure [14,](#page-21-0) there are notable systematic disagreements in each dataset. In particular, the behavior of the energy-phase curves appears to agree best at maximum accelerating phase for each  $V_s$  shown, including Tank/Buncher-3.

There are two separate contributions to this disagreement: measurement uncertainty and model error. With respect to measurement uncertainty, we note that as a result of the predicted growth in energy spread away from the optimum accelerating  $(V_s, \phi)$  means the beam distribution at the HEBT1 diagnostic station will be broad and faint. This makes the identification of an energy centroid difficult on the harp detector. For Tank-2 and Tank-3, the operator estimated energy spread at the Prague-Harp is graphed in Figure [16.](#page-23-0)

For each tank shown in Fig. [16,](#page-23-0) we observe the energy spread growing rapidly on either sides of the optimum accelerating phase. Observationally, this means two things. First, as the beam energy spread broadens, the observable distribution on the Prague-Harp monitor likewise broadens and correspondingly drops in intensity, making the observation more difficult.

Second, the HEBT1 diagnostic station has not been designed to allow for a full energy spread measurement at any phase. Consequently, as the energy spread grows, the longitudinal energy distribution exceeds the energy range of the beam monitor, meaning it is difficult to identify the precise location of the energy distribution centroid. This is done by eye and is therefore susceptible to interpretation, particularly when attempting to measure the centoid of a broad, low intensity distribution.

It is also important to note that this model uses fields  $\mathcal{E}(s)$  which have been generated by solution of an electrostatic boundary value problem, in other words not an RF simulation. Summary examination of the fields in Section [3](#page-7-0) does show a tendency for the peak heights to vary with respect to each other, moreso than is implied by the bead-pulls shown in the literature. While an improvement certainly stands to be made by generating such distributions in a full time-dependant code such as CST-MWS, the static fields in this document appear sufficient to predict the operational envelope of the machine, which always uses IH tanks at maximum accelerating phase.

With regards to Tank-3, Fig. [16](#page-23-0) suggests a band of about 60 $\degree$  where  $\Delta E/\Delta E_0$  remains near or below unity and the likely measurement error small. Beyond this, the growth in energy spread renders the measurement more uncertain, especially at higher energy. As a testament to this, note that the red curve for EPICS amplitude 400 in Fig. [16](#page-23-0) suggests a relatively small energy spread for each of the measured phases. Inspection of Figure [14](#page-21-0) (middle) for Tank-3 shows that the lower two amplitude energy-phase lines, which starts with A=400, agree most closely with the model, especially when compared to the higher energies in the set.

For Buncher-3, the disagreement is more interesting, owing to a note that was written during data acquisition, shown in Figure [17,](#page-25-1) where the effect of DTL Buncher-3 upon the beam's energy spread was sketched. The input distribution (left) exits the buncher with a double-peak feature, shown on the right. The note specifies that the high energy tail, corresponding to the narrower, higher energy distribution, was measured off-peak phase.

This can be appreciated in Figure [14,](#page-21-0) bottom-right, where for each phase-scan, the measured beam energy is seen to vary less than what TRANSOPTR predicted, on the positive side of the optimum accelerating phase. The agreement is much better to the negative side, where such deformation was not seen. Certainly, a component of the disagreement between TRANSOPTR and the actual DTL measurement must be due to the difference in  $\mathcal{E}(s)$ . Nevertheless, it would be hazardous to attempt to fine-tune  $\mathcal{E}(s)$  on a beam-based dataset, itself which relies upon a set of challenging energy measurements.



<span id="page-25-1"></span>Figure 17: Note taken by author during 2017 data acquisition for DTL Buncher-3. The HEBT1 Magnet (Prague) cooling water temperature is recorded. Observe the recorded energy distribution deformation, where the input beam distribution (left) exits the buncher with a double peak feature, shown on the right.

### <span id="page-25-0"></span>**5.2 Longitudinal Tune and Topological Ramp**

The script topology was used to perform grid scans of the entire  $(V_s, \phi)$  configuration space of each cavity, in each case spanning the allowable range for both tuning parameters. As in [\[21\]](#page-35-10), by tracking the output inverse longitudinal momentum  $P<sub>z</sub>$  at the location of the HEBT1 energy diagnostic, we obtain a signal which is resonant when the energy spread is minimized at a given output energy. Square grid dimensions of 100 or 150 were used to provide sufficient resolution, particularly for the inverted longitudinal momentum, which defines sometimes narrow paths of minimized energy spread through the configuration space.

Output topological scans for the entire DTL are shown in Figures [18,](#page-26-0) [19](#page-27-0) and [20.](#page-28-0) The inverse z-momentum is shown on the left while the energy is shown on the right, also featuring the extracted path of optimum  $P_z^{-1}$ . Each phase-voltage ramp for each cavity is bounded by green triangles, denoting the (lower) injection and (higher) design output energy, from Table [1.](#page-2-1) Note that the ISAC-DTL RF amplifiers are usually powered on at around A=400, in many cases exceeding the bunching energy, especially in the IH tanks. Observe the relatively narrow band of observable  $P<sub>z</sub>$  for Tank-3 in Fig. [20,](#page-28-0) when compared to other accelerating tanks in the machine. To the right of each plot is the output energy for each tank and buncher, also versus phase and voltage.

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<span id="page-26-0"></span>Figure 18: TRANSOPTR grid scans for (Left)  $P_z^{-1}$  and (Right) energy, for (Top) DTL Tank-1, (Middle) DTL Buncher-1 and (**Bottom**) DTL Tank-2. A/q=30/5 beam used at 0.153 MeV/u initial injection, with preceding cavities set to design energies as per Table [1.](#page-2-1) The energy ramp with minimized  $P_z^{-1}$  is shown as black dots, bounded by green triangles, which denote the configuration space location of the design injection and output energy for each tank and buncher.

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<span id="page-27-0"></span>Figure 19: TRANSOPTR grid scans for (Left)  $P_z^{-1}$  and (Right) energy, for (Top) DTL Buncher-2, (Mid**dle**) Tank-3 and (**Bottom**) DTL Buncher-3. A/q=30/5 beam used at 0.153 MeV/u initial injection, with preceding cavities set to design energies as per Table [1.](#page-2-1) The energy ramp with minimized  $P_z$ is shown as black dots, bounded by green triangles, which denote the configuration space location of the design injection and output energy for each tank and buncher.



<span id="page-28-0"></span>Figure 20: TRANSOPTR grid scans for (Left)  $P_z^{-1}$  and (Right) energy, for (Top) DTL Tank-4 and (**Bottom**) DTL Tank-5. A/q=30/5 beam used at 0.153 MeV/u initial injection, with preceding cavities set to design energies as per Table [1.](#page-2-1) The energy ramp with minimized  $P<sub>z</sub>$  is shown as black dots, bounded by green triangles, which denote the configuration space location of the design injection and output energy for each tank.

The data shows how TRANSOPTR can be used to explore the longitudinal tune of a variable energy linac, in this case with a prescribed design energy for each tank and buncher. Future investigations can also explore these design energies. The resulting energy-phase ramps are shown in Figure [21](#page-29-0) for the IH tanks with the EPICS voltage parameter as the xaxis. The bunchers are shown in Fig. [22,](#page-30-0) showing the the cosecant-like energy ramp [\[23\]](#page-35-12). These plots show the longitudinal tune of the machine (purple), for each corresponding tank output energy (red), at a given voltage setting  $V_s$  in the control system. The computed longitudinal tune of the DTL, corresponding to the energies of Table [1](#page-2-1) with minimized  $P_z$  at the position of the HEBT1 high energy diagnostic, is presented in Table [12.](#page-31-1)



<span id="page-29-0"></span>Figure 21: TRANSOPTR Computed DTL Energy-Phase ramps, against EPICS voltage parameter.



<span id="page-30-0"></span>Figure 22: TRANSOPTR Computed DTL Energy-Phase ramps, against EPICS voltage parameter. Note the pure bunching energy, where the energy ramp reverses.



<span id="page-31-1"></span>Table 12: TRANSOPTR computed longitudinal tune setpoints for tank voltage and phases, producing tank/buncher energy output as specified in Table [1.](#page-2-1)

The longitudinal  $A/q=6$  tune in Table [12](#page-31-1) was found while completely ignoring the transverse beam parameters. It is also noted that the bunchers here have been set according to the same prescription as the IH tanks, at maximum acceleration, giving a tune which underbunches the beam longitudinally, procuding a longer time-structure while still having a minimized output energy spread. This tune is intended for model development and not operational use, as of yet.

### <span id="page-31-0"></span>**5.3 Transverse ISAC-DTL Envelope**

As a demonstration of the optr-DTL's ability to produce ISAC linac tunes, the design tune exiting the MEBT section, at the location of DTL:FC0, has been fed into the ISAC-DTL and flown up to the HEBT1-Prague diagnostic station. Recalling that in Section [4,](#page-18-0) a set of optr TWISSMATCH calls were inserted as element fortline in the /acc database, automatically being included in sy.f for each xml2optr generation of the DTL sequence.

An advantage of this database inclusion of transverse focussing calls is that match parameters can easily be loaded into the database, facilitating the adjustment of the tune. Moreover, the constant presence of 8 separate TWISSMATCH lines, 2 after each triplet, means each can now be optimized via optr's internal optimizer [\[24\]](#page-35-13), by enabling the optimize flag in data.dat for each triplet. Table [13](#page-32-1) shows the Twiss parameters that were used for the present tune generation.

Each triplet has been individually optimized using this method, with the resulting transverse tune shown in Table [14.](#page-33-0) As a final demonstration of the model, an  $A/q = 30u/5$  tune, with injection energy of 0.153 MeV/u and injection parameters from Tables [11](#page-22-2) and [10,](#page-22-1) is shown in Figure [23,](#page-32-2) featuring acceleration up to 1.53 MeV/u, corresponding to Tank-5 design output energy.

TWISSMATCH-1		TWISSMATCH-2		TWISSMATCH-3		TWISSMATCH-4	
$(\alpha_x,\beta_x)$	$(\alpha_y,\beta_y)$	$(\alpha_x,\beta_x)$	$(\alpha_y,\beta_y)$	$(\alpha_x,\beta_x)$	$(\alpha_y,\beta_y)$	$(\alpha_x, \beta_x)$	$(\alpha_y,\beta_y)$
(2.52, 91.07)	(1.86, 91.72)	$(2.32, 216.44)$ $(1.42, 98.88)$ $(0.50, 60.25)$				$(1.28, 136.24)$ $(2.24, 195.97)$	(0.99,74.74)

<span id="page-32-1"></span><span id="page-32-0"></span>Table 13: Loaded DTL transverse Twiss parameters, used for the TWISSMATCH statements in dtl\_db0.xml (Sec. [4\)](#page-18-0), used to define the transverse DTL tune in the present analysis. The param-eters are derived from [\[2\]](#page-34-2). All Twiss- $\beta_{x,y}$  parameters are in [cm].



<span id="page-32-2"></span>Figure 23: TRANSOPTR 2rms ISAC-DTL envelope simulation with Tank-5 output at 1.53 MeV/u, starting at DTL:FC0 and terminating at the HEBT1-Prague Harp energy diagnostic. An  $A/q =$ 30u/5 beam is injected at 0.153 MeV/u following design parameters obtained from [\[9\]](#page-34-9). Simulation starting parameters shown in Tables [11](#page-22-2) & [10](#page-22-1) and TWISSMATCH parameters from Table [13.](#page-32-1) The model longitudinal tune is shown in Table  $12$ . Twiss parameter matching point  $s$ -coordinates are shown as red triangles.



<span id="page-33-0"></span>Table 14: TRANSOPTR computed DTL quadrupole settings for Tank-5 energy tune from Table [12.](#page-31-1) Quadrupole BI parametrization found in [\[15\]](#page-35-4).

# **6 Conclusion**

The TRANSOPTR-DTL presented herein is intended as a tool for research & development, enabling realtime simulation and study of operational tunes and further investigations of machine configuration. Starting from design drawings, tube dimensions and a handful of boundary conditions, the optr model has been built, then used to compute a full longitudinal DTL tune through model analysis. This was then verified against  $^{16}O^{4+}$  energy-phase curves, which were also used to generate a calibration between the model and the EPICS control system's voltage variables.

Good agreement between model and machine was found with all cavities, except Tank-3 and Buncher-3. Consideration of the challenging nature of the energy measurement for Tank-3 may at least partially explain this disagreement. For Buncher-3 it was noted that during acquisition the high-energy peak, not the broad body of the energy distribution, was tracked beyond the optimum accelerating phase, producing a systematic disagreement with the optr prediction. While these factors possibly explain some of the discrepancies, from Fig. [14,](#page-21-0) it is evident that more scrutiny is called for, particularly when considering the degree of agreement for the remainder of the linac. An investigation into the longitudinal dynamics using a multiparticle model in the code LORASR is also underway.

The availability of a full TRANSOPTR end-to-end model for the ISAC-I linac now opens the door to the development of model-coupled accelerator tuning software and techniques, by exploiting the TRIUMF-HLA infrastructure. To this end, the transverse DTL tune will also see more scrutiny.

# <span id="page-34-0"></span>**7 Acknowledgements**

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# <span id="page-36-0"></span>**Appendices**

# <span id="page-36-1"></span>**A DTL Energy-Phase Mapping Original Run-Plan**

## <span id="page-36-2"></span>**A.1 Run Plan**

A systematic approach will ensure consistent data acquisition practices. The OLIS microwave source will produce several nanoamps of 16O+, stripped to 4+ in MEBT. Beam will be tuned from OLIS through the RFQ and MEBT section and initially drifted through the DTL, to be measured on the Prague magnet at the expected 0.153 MeV/u drift energy. Prior to DTL energy phasing investigations, a parasitic test of the RFQ's energy acceptance and output will be performed, providing data for ongoing RFQ modellization efforts. An OLIS emittance scan will also be performed for beam characterization. Following this, the energy-phase investigation of the DTL will carry forward.

This will be accomplished by incrementing DTL tank amplitudes by a fixed amount finding the corresponding maximum accelerating phase. These are noted, then the phase will be scanned on either side of maximum acceleration at constant intervals and the corresponding output energy profile centroid and FWHM recorded. The measurements are to be carried out on each of the five DTL accelerating tanks and its 3 bunchers, providing a mapping of the whole linac.

### <span id="page-36-3"></span>**A.2 Protocol**

All required RPM scans should both be graphically saved to the ISAC e-log, and saved as a data file for later analysis. All DTL energy-phase measurements will be logged in a provided run logbook, written in ink.

### <span id="page-36-4"></span>**A.2.1 OLIS Emittance Measurement**

- 1. Tune beam up to the Prague magnet, drifting through the DTL, with the MEBT rebuncher off.
- 2. Perform an OLIS emittance measurement.

### <span id="page-37-0"></span>**A.2.2 Parasitic RFQ Experiment**

Validation of PARMTEQ and TRANSOPTR model performance requires a measurment of the RFQ output transmission and energy for varying input beam energy.

Once beam is tuned through the RFQ following standard practices, the OLIS voltage will be scanned around the input energy. Both the RFQ transmission between ILT:FC49 and MEBT:FC5 will be recorded, in addition to the FWHM and energy centroid on the Prague magnet.

- 1. Find both the lower and upper source biases that completely kill transmission through the RFQ, as measured on MEBT:FC5.
- 2. Go to the lower voltage.
- 3. Increment the source bias by a value equal to one tenth the difference between the low and high injection voltages.
- 4. For each step, record the RFQ transmission, profile on MEBT:FC5, time profile on MEBT:FFC5 and energy profile on the Prague HARP.

### <span id="page-37-1"></span>**A.2.3 DTL Energy-Phase Mapping**

- 1. Ensure DTL drifting tune is fully optimized using standard practices.
- 2. Document initial drifting tune through DTL. Scan ILT:RPM37, MEBT:RPM5, HEBT:RPM5, take transmission snapshot and record Prague HARP and beam energy.
- 3. Turn on DTL Tank1 at EPICS amplitude 400.
- 4. Reoptimize DTL Quadrupoles, find beam on Prague magnet.
- 5. Find maximum accelerating phase, take transmission snapshot.
- 6. With a step size of 10°, record the output energy and FWHM for 10 points below and 10 points above the maximum accelerating phase.
- 7. Increase Tank amplitude by 50 EPICS units.
- 8. Repeat from Step 4 until at maximum allowable tank energy.
- 9. **The last amplitude step should be set to exactly the maximum allowable amplitude.**
- 10. The above will also be repeated for DTL bunchers.

## <span id="page-38-0"></span>**B** TRANSOPTR**-DTL** sy.f

```
SUBROUTINE TSYSTEM
 COMMON/SCPARM/QSC,ISC,CMPS
 COMMON/MOM/P,BRHO,pMASS,ENERGK,GSQ,ENERGKi,charge,current
 COMMON/BLOC1/RFA1,RFP1,QM1,QM2,QM3,RFA2,RFP2,RFA3,RFP3,QM4,QM5,QM6
&,RFA4,RFP4,RFA5,RFP5,QM7,QM8,QM9,RFA6,RFP6,RFA7,RFP7,QM10,QM11,QM1
&2,RFA8,RFP8,QM13,QM14,QM15,QM16,QM17,QM18,QM19
 CMPS=1.74 ! Number of cm per step, for plotting only
 wo=1.0 ! Weight aberration from optical elements
 ! start-dtl
 call drift(0.83,".")
call marker('DTL:FC0')
call drift(-0.83,".")
! RF cavity/linac: ISAC1:DTL1
call linacn(100,1001,6.79376e-06*RFA1**2.0 - 0.00219132*RFA1 +
&2.55377,32.746,1.0608e+08,1.0*RFP1,'ISAC1:DTL1')
call drift(2.25875,".")
 ! Magnetic quadrupole DTL:Q1
 call mquad(-213.4777*tanh(1.054896e-24*QM1**5.0 +
&3.599039e-15*QM1**3.0 + 2.210222e-05*QM1),1.1988,6.5,wo,'DTL:Q1')call drift(4.99845,".")
 ! Magnetic quadrupole DTL:Q2
call mquad(251.3987*tanh(5.982539e-25*QM2**5.0 +
&2.560897e-15*QM2**3.0 + 1.9732e-05*QM2),1.1988,9.5,wo,'DTL:Q2')
call drift(4.95,".")
 ! DTL:XCB2
! DTL:YCB2
call drift(0.04845,".")
 ! Magnetic quadrupole DTL:Q3
call mquad(-213.4777*tanh(1.054896e-24*QM3**5.0 +
&3.599039e-15*QM3**3.0 + 2.210222e-05*QM3),1.1988,6.5,wo,'DTL:Q3')
call drift(1.7,".")
call TWISSMATCH( 1, 2.52,91.07, 0.5, 1)
 call TWISSMATCH( 3, 1.86,91.72, 0.5, 1)
call drift(-0.004235,".")
! RF cavity/linac: ISAC1:BUNCH1
 call linacn(101,1001,0.007986697*RFA2 +
&0.3061407,13.050,1.0608e+08,1.0*RFP2,'ISAC1:BUNCH1')
call drift(4.85129,".")
! RF cavity/linac: ISAC1:DTL2
 call linacn(102,1001,0.006828576*RFA3 -
&0.3663625,54.775,1.0608e+08,1.0*RFP3,'ISAC1:DTL2')
```

```
call drift(2.19065,".")
 ! Magnetic quadrupole DTL:Q4
 call mquad(213.4777*tanh(1.054896e-24*QM4**5.0 +
&3.599039e-15*QM4**3.0 + 2.210222e-05*QM4),1.1988,6.5,wo,'DTL:Q4')
 call drift(4.9985,".")
 ! Magnetic quadrupole DTL:Q5
 call mquad(-251.3987*tanh(5.982539e-25*QM5**5.0 +
&2.560897e-15*QM5**3.0 + 1.9732e-05*QM5),1.1988,9.5,wo,'DTL:Q5')
 call drift(4.95,".")
 ! DTL:XCB5
! DTL:YCB5
call drift(0.0484,".")
 ! Magnetic quadrupole DTL:Q6
 call mquad(213.4777*tanh(1.054896e-24*QM6**5.0 +
&3.599039e-15*QM6**3.0 + 2.210222e-05*QM6), 1.1988, 6.5, wo, 'DTL:Q6')call drift(1.7,".")
call TWISSMATCH( 1, 2.32,216.44, 0.5, 1)
call TWISSMATCH( 3, 1.42,98.88, 0.5, 1)
call drift(-0.002055,".")
! RF cavity/linac: ISAC1:BUNCH2
 call linacn(103,1001,-1.5917e-06*RFA4**2.0 + 0.008996733*RFA4 -
&0.718549,14.850,1.0608e+08,1.0*RFP4,'ISAC1:BUNCH2')
call drift(4.84989,".")
 ! RF cavity/linac: ISAC1:DTL3
 call linacn(104,1001,0.007007824*RFA5 -
&0.4691255,81.775,1.0608e+08,1.0*RFP5,'ISAC1:DTL3')
call drift(2.19195,".")
 ! Magnetic quadrupole DTL:Q7
 call mquad(-213.4777*tanh(1.054896e-24*QM7**5.0 +
&3.599039e-15*QM7**3.0 + 2.210222e-05*QM7),1.1988,6.5,wo,'DTL:Q7')
call drift(4.9984,".")
 ! Magnetic quadrupole DTL:Q8
 call mquad(251.3987*tanh(5.982539e-25*QM8**5.0 +
&2.560897e-15*QM8**3.0 + 1.9732e-05*QM8),1.1988,9.5,wo,'DTL:Q8')
call drift(4.95,".")
 ! DTL:XCB8
 ! DTL:YCB8
call drift(0.0485,".")
 ! Magnetic quadrupole DTL:Q9
call mquad(-213.4777*tanh(1.054896e-24*QM9**5.0 +
&3.599039e-15*QM9**3.0 + 2.210222e-05*QM9),1.1988,6.5,wo,'DTL:Q9')
call drift(1.7,".")
 call TWISSMATCH( 1, 0.50,60.25, 0.5, 1)
call TWISSMATCH( 3, 1.28,136.24, 0.5, 1)
 call drift(-0.00209,".")
 ! RF cavity/linac: ISAC1:BUNCH3
```

```
call linacn(105,1001,0.006496419*RFA6 +
&0.1508214,17.349,1.0608e+08,1.0*RFP6,'ISAC1:BUNCH3')
 call drift(4.8499,".")
 ! RF cavity/linac: ISAC1:DTL4
call linacn(106,1001,0.004845639*RFA7 -
&0.1695915,94.775,1.0608e+08,1.0*RFP7,'ISAC1:DTL4')
 call drift(2.19189,".")
 ! Magnetic quadrupole DTL:Q10
call mquad(213.4777*tanh(1.054896e-24*QM10**5.0 +
&3.599039e-15*QM10**3.0 +
&2.210222e-05*QM10),1.1988,6.5,wo,'DTL:Q10')
call drift(4.9985,".")
 ! Magnetic quadrupole DTL:Q11
 call mquad(-251.3987*tanh(5.982539e-25*QM11**5.0 +
&2.560897e-15*QM11**3.0 + 1.9732e-05*QM11),1.1988,9.5,wo,'DTL:Q11')
 call drift(4.95,".")
 ! DTL:XCB11
! DTL:YCB11
call drift(0.0484,".")
 ! Magnetic quadrupole DTL:Q12
call mquad(213.4777*tanh(1.054896e-24*QM12**5.0 +
&3.599039e-15*QM12**3.0 +
&2.210222e-05*QM12),1.1988,6.5,wo,'DTL:Q12')
call drift(8.8656,".")
call TWISSMATCH( 1, 2.24,195.97, 0.5, 1)
call TWISSMATCH( 3, 0.99,74.74, 0.5, 1)
! RF cavity/linac: ISAC1:DTL5
call linacn(107,1001,1e-06*RFA8,102.77,1.0608e+08,1.0*RFP8,'ISAC1:
&DTL5')
! endOf_dtl_dbO
call drift(6.4599,".")
! startOf_t3d_tune_hebt
call drift(5.7751,".")
 ! HEBT:IV0
call drift(9.785,".")
call marker('HEBT:RPM0')
call drift(8.262,".")
call marker('HEBT:FC0')
call drift(20.678,".")
! HEBT:XCB0
! HEBT:YCB0
call drift(11.5,".")
 ! Magnetic quadrupole HEBT:Q1
 call mquad(-0.009228503*QM13,2.6,18.0,wo,'HEBT:Q1')
 call drift(19.0,".")
 ! Magnetic quadrupole HEBT:Q2
```

```
call mquad(0.009228503*QM14,2.6,18.0,wo,'HEBT:Q2')
call drift(72.977,".")
! HEBT:XCB2
! HEBT:YCB2
call drift(11.523,".")
! Magnetic quadrupole HEBT:Q3
call mquad(-0.009228503*QM15,2.6,18.0,wo,'HEBT:Q3')
call drift(56.0,".")
! Magnetic quadrupole HEBT:Q5
call mquad(0.009228503*QM16,2.6,18.0,wo,'HEBT:Q5')
call drift(69.0,".")
call marker('HEBT:RPM5')
! HEBT:STRP5
call drift(10.54,".")
call marker('HEBT:FC5')
! HEBT:SCD5
call drift(15.129,".")
! HEBT:PSID5
call drift(12.914,".")
! HEBT:IV8
call drift(37.274,".")
! HEBT:XCB5
! HEBT:YCB5
call drift(14.143,".")
! Magnetic quadrupole HEBT:Q6
call mquad(-0.009283983*QM17,2.6,18.0,wo,'HEBT:Q6')
call drift(20.0,".")
! Magnetic quadrupole HEBT:Q7
call mquad(0.009477347*QM18,2.6,32.5,wo,'HEBT:Q7')
call drift(20.0,".")
! Magnetic quadrupole HEBT:Q8
call mquad(-0.009228503*QM19,2.6,18.0,wo,'HEBT:Q8')
call drift(13.46,".")
! HEBT:XCB8
! HEBT:YCB8
call drift(10.08,".")
! HEBT:XCOL8A
call drift(9.8153,".")
! endOf_hebt_db0
call drift(167.162,".")
! Magnetic dipole HEBT1:MB0
call edge(0.0,153.719,90.0,0.0,0.0,0.0,5.0,0.0,wo)
call bend(153.719,90.0,0.0,'HEBT1:MB0')
call edge(0.0,153.719,90.0,0.0,0.0,0.0,5.0,0.0,wo)
call drift(86.143,".")
! HEBT1:IV0
```

```
call drift(67.3875,".")
call marker('HEBT1:HARP0')
call print_transfer_matrix
return
end
```
# <span id="page-42-0"></span>**C** TRANSOPTR**-DTL** data.dat

134.6 0.0 250.0 0 ! DTL:Q12:CUR QM12 A

0.0 0.0 2e+08 0! ISAC1:DTL5:AMP:SETPT RFA8 V

4.59 0.0 0.0 27944.8 5.0 0.0 ! En[MeV], mom., brho, mass[MeV], charge, beam cur. or bun.chg. -1 5 0.01 0.0001 ! iprint, IVOPT (4/5: 4/6-D space-charge), initial RK step, RK error tol. 0 0.0 1.0 0.0 ! for external Bs (0=disable), s offset, unit of s (1=cm), unit of Bs (1=kG) 0.358748 0.0236304 0.34 0.0225397 0.1345 0.025171 ! bunch dim: x,x',y,y',z(bun. len.),dp/p 1.0 1.0 1.0 1.0 1.0 1.0 1.0 1.0 ! 1 means x,y,z in cm, x',y',dp/p in rad (dimensionless) 3 1 2 -0.887894 3 4 -0.896195 5 6 -0.62 35 757.58 0.0 2500.0 0 ! ISAC1:DTL1:AMP:SETPT RFA1 V 10.0 -360.0 360.0 0 ! ISAC1:DTL1:PHASE:SETPT RFP1 deg 129.1 0.0 250.0 0 ! DTL:Q1:CUR QM1 A 111.1 0.0 250.0 0 ! DTL:Q2:CUR QM2 A 99.26 0.0 250.0 0 ! DTL:Q3:CUR QM3 A 487.48 -38.33 2500.0 0 ! ISAC1:BUNCH1:AMP:SETPT RFA2 V 187.34 -360.0 360.0 0 ! ISAC1:BUNCH1:PHASE:SETPT RFP2 deg 1051.28 0.0 2500.0 0 ! ISAC1:DTL2:AMP:SETPT RFA3 V 75.14 -360.0 360.0 0 ! ISAC1:DTL2:PHASE:SETPT RFP3 deg 109.7 0.0 250.0 0 ! DTL:Q4:CUR QM4 A 119.5 0.0 250.0 0 ! DTL:Q5:CUR QM5 A 102.9 0.0 250.0 0 ! DTL:Q6:CUR QM6 A 772.15 0.0 2500.0 0 ! ISAC1:BUNCH2:AMP:SETPT RFA4 V 214.53 -360.0 360.0 0 ! ISAC1:BUNCH2:PHASE:SETPT RFP4 deg 1055.7 0.0 2500.0 0 ! ISAC1:DTL3:AMP:SETPT RFA5 V 257.84 -360.0 360.0 0 ! ISAC1:DTL3:PHASE:SETPT RFP5 deg 134.9 0.0 250.0 0 ! DTL:Q7:CUR QM7 A 147.3 0.0 250.0 0 ! DTL:Q8:CUR QM8 A 112.8 0.0 250.0 0 ! DTL:Q9:CUR QM9 A 613.58 -23.22 2500.0 0 ! ISAC1:BUNCH3:AMP:SETPT RFA6 V 124.17 -360.0 360.0 0 ! ISAC1:BUNCH3:PHASE:SETPT RFP6 deg 1277.35 0.0 2500.0 0 ! ISAC1:DTL4:AMP:SETPT RFA7 V 303.33 -360.0 360.0 0 ! ISAC1:DTL4:PHASE:SETPT RFP7 deg 121.3 0.0 250.0 0 ! DTL:Q10:CUR QM10 A 156.4 0.0 250.0 0 ! DTL:Q11:CUR QM11 A

